

## HIDDEN REGULARITY FOR A STRONGLY NON LINEAR

### WAVE EQUATION

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#### ABSTRACT

In this paper we consider the nonlinear wave equation,

$$\begin{cases} u'' - \Delta u + f(u) = V & \text{in } Q = \Omega \times ]0, T[; \\ u(0) = u_0, u'(0) = u_1 & \text{in } \Omega, \\ u(x, t) = 0 & \text{on } \Sigma = \Gamma \times ]0, T[ \end{cases} \quad (*)$$

where  $f$  is a continuous function satisfying

$$\limsup_{|s| \rightarrow +\infty} \frac{f(s)}{s} > -\infty \quad (**)$$

and  $\Omega$  is a bounded domain of  $\mathbb{R}^n$  with smooth boundary  $\Gamma$ . We prove that

there exist a solution for (\*) that satisfies the regularity conditions:  $\frac{\partial u}{\partial \eta} \in L^2(\Sigma)$ .

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Moreover we have that there exist a constant  $C > 0$  such that,

$$\left| \frac{\partial u}{\partial \eta} \right| \leq C \{ E(0) + |V|_Q^2 \} \quad (***)$$

## 1. INTRODUCTION

Let  $\Omega$  be an open bounded set of  $\mathbb{R}^n$ , with boundary  $\Gamma$  of class  $C^2$ . Set  $Q = \Omega \times ]0, T[$  and  $\Sigma = \Gamma \times ]0, T[$ . We will denote by  $(\cdot, \cdot)_\Omega$  and  $(\cdot, \cdot)_Q$  the inner product of  $L^2(\Omega)$  and  $L^2(Q)$  respectively and by  $|\cdot|_\Omega$ ,  $|\cdot|_Q$  and  $\|\cdot\|$ , the norms in  $L^2(\Omega)$ ,  $L^2(Q)$ , and  $H_0^1(\Omega)$  respectively. We consider the nonlinear problem:

$$\begin{cases} u'' - \Delta u + f(u) = V & \text{in } Q; \\ u(0) = u_0, u'(0) = u_1 & \text{in } \Omega; \\ u(x, t) = 0 & \text{on } \Sigma \end{cases} \quad (1.1)$$

In J.L.Lions [1] was study the hidden regularity for system (1.1) when  $f(s) = s^3$  and more generality for a  $f(s) = s|s|^e$ , where  $e \geq 0$ . In this work we are going to study the hidden regularity for the solution of the problem (1.1) when  $f$  is a

continuous function satisfying,

$$\limsup_{|s| \rightarrow +\infty} \frac{f(s)}{s} > -\infty \quad (1.2)$$

That is, we will show that there exist a solution  $u$  of the above problem such that the normal derivate of  $u$  belongs to  $L^2(\Sigma)$ . Moreover we will prove that there exist a constant  $C > 0$  such that:

$$\left| \frac{\partial u}{\partial \eta} \right|_{\Sigma} \leq CE_0 \quad (1.3)$$

where  $E_0$  is the initial energy of the system (1.1).

$$\left( E_0 = \frac{1}{2} |u_1|_{\Omega}^2 + \int_{\Omega} G(u_0) dx \right)$$

where  $G(s) = \int_0^s f(\eta) d\eta$ .

## 2. EXISTENCE AND HIDDEN REGULARITY

First of all we are going to construct a sequence of real numbers  $(s_\nu)_{\nu \in \mathbb{N}}$  and  $(s_{-\nu})_{\nu \in \mathbb{N}}$  satisfying the following conditions,

$$s_\nu \geq \nu, \forall \nu \in \mathbb{N}, |f(s_\nu)| \leq C + |f(s)|, \forall s \geq \nu \quad (2.1)$$

$$s_{-\nu} \geq -\nu, \forall \nu \in \mathbb{N}, |f(s_{-\nu})| \leq C + |f(s)|, \forall s \geq -\nu \quad (2.2)$$

This sequences are going to play an important role in the sequel.

**LEMMA 2.1.-** Let's  $f$  be a continuous function defined in  $\mathbb{R}$ , then there exists a sequence of real numbers,  $(s_\nu)_{\nu \in \mathbb{N}}$ , and a positive constant  $C$ , depending of  $\nu$ , satisfying condition (2.1) and (2.2).

**PROOF.-** Let's consider the following problem.

$$I_\nu = \inf \{|f(s)|; s \geq \nu\} \quad (2.3)$$

If for all  $\nu \in \mathbb{N}$ , there exist  $s_\nu \geq \nu$  such that  $f(s_\nu) = I_\nu$ , then this sequence satisfies condition (2.1). Now we can suppose that there exist a  $v_0$  such that,

$$f(s) > I_{v_0} \quad \text{for all } s \geq v_0$$

This relations imply that  $I_\nu = I_{\nu_0}$  for all  $\nu \geq \nu_0$ . Let us put  $I_0 = I_{\nu_0}$ . Since  $I_0 = \inf |s(s)|; s \geq V_0$ , there exists a sequence  $(t_k)_{k \in \mathbb{N}}$  such that,

$$f(\bar{t}_k) \rightarrow I_0 \quad (2.4)$$

from the continuity of  $f$  we conclude that  $t_k$  is not bounded, then there exist a subsequence  $(t_{k_\nu})_{\nu \in \mathbb{N}}$  satisfying:

$$t_{k_\nu} \geq V, \quad \nu \in \mathbb{N} \quad (2.5)$$

Let us put  $s_\nu = t_{k_\nu}$ , from (2.4) we obtain that there exist a constant  $\mathcal{C}$  (independent of  $\nu$ ) such that  $|f(s_\nu)| = |f(t_{k_\nu})| \leq \mathcal{C}$ , finally from (2.5) we conclude that  $(s_\nu)_{\nu \in \mathbb{N}}$  satisfies condition (2.1). By the same arguments we can prove the existence of a sequence  $(s_\nu)_{\nu \in \mathbb{N}}$  satisfying condition (2.2), only consider the problem

$$I_{-\nu} = \inf |f(s)|; s \leq -\nu,$$

and the result follows.

With the sequences  $(s_\nu)_{\nu \in \mathbb{N}}$  and  $(s_{-\nu})_{\nu \in \mathbb{N}}$  constructed in Lemma 2.1 we define

a sequence  $(f_\nu)_{\nu \in \mathbb{N}}$  of continuous function in the following way:

$$f_\nu(s) = \begin{cases} f(s) & \text{if } s_\nu \leq s \leq s_\nu; \\ f(s_\nu) & \text{if } s \geq s_\nu; \\ f(s_{-\nu}) & \text{if } s_\nu \leq s_{-\nu}. \end{cases} \quad (2.6)$$

As a consequence of Lemma 2.1 we have that the sequence  $(f_\nu)_{\nu \in \mathbb{N}}$  satisfies the following properties:

$$|f_\nu(s)| \leq C + |f(s)| \quad \text{for all } \nu \quad (2.7)$$

$$f_\nu \rightarrow f \quad \text{uniformly on bounded sets} \quad (2.8)$$

**LEMMA 2.2.-** Let's  $f$  be a continuous function satisfying condition (1.2), and  $(f_\nu)_{\nu \in \mathbb{N}}$ , the sequence defined in (2.6). Then there exist a positive constant  $C_0$  such that.

$$sf_\nu(s) \geq -C_0(s^2 + 1) \quad \forall s \in \mathbb{R} \quad \forall \nu \geq \nu_0 \quad (2.9)$$

$$\int_0^t f_\nu(s) ds \geq -C_0(t^2 + 2|t|) \quad (2.10)$$

$$\left| \int_0^t f_\nu(s) ds \right| \leq \frac{1}{2} C_0 |t(t+3)| + \int_0^t f(s) ds \quad (2.11)$$

**PROOF.-** First of we are going to prove that there exists a positive constant  $C_0$ , such that

$$f(s) \geq -C_0 s \quad \forall s \geq N \quad \text{and} \quad f(s) \leq -C_0 s \quad \forall s \leq N \quad (2.12)$$

In fact, if  $\liminf s^{-1}f(s) = +\infty$  the expression (2.12) is valid. Now we can suppose that  $\liminf s^{-1}f(s) = x < +\infty$ , then for  $\mathcal{E} > 0$ , there exist  $N > 0$  such that  $s^{-1}f(s) > x - \mathcal{E}$ , for  $|s| \geq N$ . Let us take  $C = \sup \{|f(s)|; |s| \leq N\}$ ,  $C_2 = \sup \{|sf(s)|; |s| \leq N\}$ , and put  $C_0 = \max \{C, C_1, C_2, |x - \mathcal{E}|\}$  where  $C$  is the constant in (2.7), certainly for this  $C_0$ , condition (2.12) is valid. Finally multiplying the relations in (2.12) by  $s$  ( $|s| \geq N$ ), we have by the definition of  $C_0$ , that the first part of (2.9) is valid. The second part of (2.9) follows from (2.1), (2.2), (2.6) and also, the definition of  $C_0$ , for  $v_0 = N$ .

In order to prove (2.11), let us note that from (2.12) follows that:

$$f_\nu(s) \geq -C_0(s+1) \quad \forall s \geq 0 \quad \text{and} \quad f_\nu(s) \leq -C_0(s-1) \quad \forall s \leq 0 \quad (2.13)$$

Integrating this expression we obtain (2.10). In order to obtain (2.11) let us consider relation (2.7) then we have:

$$\left| \int_0^t f_\nu(s) ds \right| \leq \int_0^t |f_\nu(s)| ds C |t| + \int_0^t |f(s) ds| \quad \forall t \in \mathbb{R} \quad (2.14)$$

From (2.12) we obtain that  $f(s) \geq -C_0(s+1) \quad \forall s \geq 0$ , which imply that  $|f(s)| \leq f(s) + 2C_0(s+1) \quad \forall s \geq 0$ . Then we have:

$$\int_0^t |f(s) ds| \leq \int_0^t f(s) ds + C_0 t(t+2) \quad \forall t \geq 0 \quad (2.15)$$

and since  $f(s) \leq |f(s)|$  we obtain,

$$\int_0^t |f(s)| ds \leq \int_0^t f(s) ds \quad \forall t \geq 0 \quad (2.16)$$

Finally from (2.15) and (2.16) we obtain (2.11).

Let us denote by  $G_\nu(t) = \int_0^t f_\nu(s) ds$ , then we have that

$$G_\nu \rightarrow G \quad \text{uniformly on bounded sets.} \quad (2.17)$$

Before to prove the main result of this paper we will prove an identity that



will be fundamental in that follows.

**LEMMA 2.3.-** Let  $h$  be a continuous function. Let  $q = (q_k)$  a field of vectors of class  $[C^1(\bar{\Omega})]^n$ . Then for all  $W$  satisfying

$$W \in L^2(0, T; H_0^1(\Omega) \cap H^2(\Omega)), h(W) \in L^1(Q) \quad (2.18)$$

$$W' \in L^2(0, T; H_0^1(\Omega)) \quad (2.19)$$

$$W'' \in L^2(0, T; L^2(\Omega)) \quad (2.20)$$

$$\begin{cases} W'' - \Delta W + h(W) = V & \text{in } Q, \\ W(0) = W_0, W'(0) = W_1 & \text{in } \Omega; \\ W(x, t) = 0 & \text{on } \Sigma. \end{cases} \quad (2.21)$$

Where  $H$  is the primitive of  $h$ . Then we have.

$$\begin{aligned} \frac{1}{2} \int_{\Sigma} q_k n_k \left| \frac{\partial W}{\partial \eta} \right|^2 d\Sigma &= \left[ (W'(t), q_k \frac{\partial W(t)}{\partial x_k})_0^T + \right. \\ &+ \int_Q \frac{\partial q_k}{\partial x_j} \left\{ |W'|^2 - |\nabla W|^2 - 2H(W) \right\} dx dt \\ &+ \int_Q \frac{\partial q_k}{\partial x_j} \times \frac{\partial W}{\partial x_k} \times \frac{\partial W}{\partial x_j} dx dt - \int_Q V q_k \frac{\partial W}{\partial x_k} \end{aligned}$$

**PROOF.**-Let us multiply (2.21), by  $q_k \frac{\partial W}{\partial x_k}$ , then we have that:

$$\begin{aligned} \int_Q \{W'' - \Delta W + h(W)\} q_k \frac{\partial W}{\partial x_k} dx dt &= \\ \int_Q V q_k \frac{\partial W}{\partial x_k} dx dt & \end{aligned} \quad (2.22)$$

Let us denote by:

$$\begin{aligned} I_1 &= \int_Q W'' q_k \frac{\partial W}{\partial x_k} dx dt \\ I_2 &= \int_Q \Delta W q_k \frac{\partial W}{\partial x_k} \end{aligned}$$

then we have:

$$\begin{aligned} I_1 &= \left[ (W'(t), q_k \frac{\partial W}{\partial x_k}(t)) \right]_0^T - \int_0^T \int_Q W' q_k \frac{\partial W'}{\partial x_k} dx dt \\ &= \left[ (W'(t), q_k \frac{\partial W(t)}{\partial x_k}) \right]_0^T - \frac{1}{2} \int_0^T \int_Q q_k \frac{\partial |W'|^2}{\partial x_k} dx dt \end{aligned}$$

from where we obtain that:

$$I_1 = \left[ \left( W'(t), q_k \frac{\partial W(t)}{\partial x_k} \right)_{\Omega} \right]_0^T + \frac{1}{2} \int_0^T \frac{\partial q_k}{\partial x_k} |W'|^2 dx dt \quad (2.23)$$

on the other hand we have that:

$$\begin{aligned} I_2 &= - \int_Q \frac{\partial W}{\partial x_j} x \frac{\partial}{\partial x_j} \left\{ q_k \frac{\partial W}{\partial x_k} \right\} dx dt + \int_{\Sigma} q_k \frac{\partial W}{\partial x_k} x \frac{\partial W}{\partial \eta} d\Sigma \\ &= - \int_Q \frac{\partial W}{\partial x_k} \frac{\partial q_k}{\partial x_k} x \frac{\partial W}{\partial x_k} dx dt - \int_Q \frac{\partial W}{\partial x_j} x \frac{\partial^2 W}{\partial x_k \partial x_j} x q_k dx dt. \\ &\quad + \int_{\Sigma} q_k \frac{\partial W}{\partial x_k} x \frac{\partial W}{\partial \eta} d\Sigma. \\ &= - \int_Q \frac{\partial W}{\partial x_k} x \frac{\partial q_k}{\partial W_j} x \frac{\partial W}{\partial x_k} dx dt - \frac{1}{2} \int_Q q_k \frac{\partial}{\partial x_k} |\nabla W|^2 dx dt + \int_{\Sigma} q_k \frac{\partial W}{\partial x_k} x \frac{\partial W}{\partial \eta} d\Sigma. \end{aligned}$$

But since  $W = 0$  on  $\Sigma$  we have that:

$$\frac{\partial W}{\partial x_k} = \eta_k \frac{\partial W}{\partial \eta} \text{ on } \Sigma.$$

and

$$|\nabla W|^2 = \left| \frac{\partial W}{\partial \eta} \right|^2 \text{ on } \Sigma.$$

we have:

$$\begin{aligned}
 I_2 = & - \int_Q \frac{\partial W}{\partial x_j} x \frac{\partial W}{\partial x_k} x \frac{\partial q_k}{\partial x_j} dx dt + \frac{1}{2} \int_Q \frac{\partial q_k}{\partial x_k} |\nabla W|^2 dx dt. \\
 & + \frac{1}{2} \int_{\Sigma} q_k \eta_k \left| \frac{\partial q_k}{\partial x_k} \right|^2 d\Sigma
 \end{aligned} \tag{2.24}$$

finally since

$$\int_Q h(W) q_k \frac{\partial W}{\partial x_k} dx dt = \int_Q \frac{\partial}{\partial x_k} H(W) q_k = \int_Q H(W) \frac{\partial}{\partial x_k} q_k. \tag{2.25}$$

we have from (2.22), (2.23), (2.24) and (2.25) that the result follows.

**REMARK 2.4.-** From Lemma 2.3 taking  $q = (q_k)$  a field of vectors of class  $[C^1(\bar{\Omega})]^n$ , such that

$$q_k = \eta_k \text{ on } \Sigma.$$

and putting:

$$\mathcal{C} = \sup \left\{ |q_k(x)|, \left| \frac{\partial q_k}{\partial x_k}(x) \right|; k, j = 1, \dots, n \text{ and } x \in \bar{\Omega} \right\}$$

we have:

$$\begin{aligned} \frac{1}{2} \int_{\Sigma} \left| \frac{\partial W}{\partial \eta} \right|^3 d\Sigma &\leq 2n \sup_{[0,T]} J(t) + \\ + CnT \sup_{[0,T]} J(t) + Cn \int_Q |H(W)| dx dt + \\ + 2C (\sup J(t)) + \frac{Cn}{2} |V|_Q^2 + CT \sup J(t) \end{aligned}$$

From where we have:

$$\begin{aligned} \frac{1}{2} \int_{\Sigma} \left| \frac{\partial W}{\partial \eta} \right| d\Sigma &\leq C(n+1)(2+T) \sup_{[0,T]} J(t) + \\ + Cn \frac{1}{2} |V|_Q^2 + \int_Q |H(W)| dx dt \end{aligned} \tag{2.26}$$

REMARK 2.5.- From (2.21) we have that

$$W'' - \Delta W + h(W) + bW = V + bW$$

mutiplying this expression by  $W'$  and integrating in  $\Omega$  we have:

$$\frac{d}{dt} \left\{ J(t) + \int_Q H(W) + b|W|^2 dx \right\} = (V, W')_{\Omega} + b(V, W')_{\Omega}$$

where  $J(t) = \frac{1}{2} \{ |W'(t)|^2 + \|W(t)\|^2 \}$ . If we put  $C_0 = \max \{1, b+1, bc^3\}$

(where  $c$  is such that  $|\cdot|_{\Omega} \leq c \|\cdot\|$ ) we obtain, after integrate from 0 to  $t$ , that:

$$J(t) + \int_Q \{H(W) + b(W)^2\} dx \leq \frac{1}{2} |V|_Q^2 + 2C_0 E(0) + C_0 \int_0^t J(s) ds \quad (2.27)$$

where  $E(t)$  is the energy associate with system (2.21), that is:

$$E(t) = J(t) + \int_{\Omega} H(W(x, t)) dx$$

**REMARK 2.6.-** Multiplying (2.21)<sub>1</sub> by  $W$ , integrating in  $Q$ , and applying Green's formul, we have that:

$$\begin{aligned} \int_Q W h(W) dx dt &= \int_Q W V dx dt + \int_0^T |W'(t)|_{\Omega}^2 dt - \\ &\quad - \int_0^T \|W(t)\|^2 dt - (W'(t), W(t))_0^T \end{aligned}$$

from where we have that:

$$\int_Q W h(W) dx dt \leq \frac{1}{2} |V|_Q^2 + (3T + 2C) \sup_{[0, T]} J(t) \quad (2.28)$$

(where  $C$  is such that  $|\cdot|_{\Omega} \leq C \|\cdot\|$ ).

Now we are condition to prove the main result of this paper:

**THEOREM 2.7.-** Let  $(u_0, u_1, V)$  be an elemente of the space

$H_0^1(\Omega) \times L^2(\Omega) \times L^2(Q)$ , and let  $f$  be a continuous function such that  $G(u_0) \in L^1(Q)$ . Then there exist a function  $u : Q \rightarrow \mathbb{R}$  satisfying

$$u \in L^\infty(0, T; H_0^1(\Omega)), u' \in L^\infty(0, T; H^2(\Omega)) \quad (2.29)$$

$$\begin{cases} u'' - \Delta u + f(u) = V & \text{in } Q; \\ u(0) = u_0, u'(0) = u_1 & \text{in } \Omega; \\ u(x, t) = 0 & \text{on } \Sigma \end{cases} \quad (2.30)$$

**REMARK 2.8.-** We are proving here that, exist one solution satisfying the last two conditions. We don't know if all solution of (1.1) satisfies this regularity result. This is an open question.

**PROOF OF THEOREM 2.7.-** Let  $(\rho_\mu)_{\mu \in \mathbb{N}}$  be a regularizant sequence on  $\mathbb{R}$ .

That is:  $\rho_\mu \in C^\infty(\mathbb{R})$ ,  $\forall \mu \in \mathbb{N}$  and:

$$\rho_\mu(s) \geq 0 \quad \forall s \in \mathbb{R} \quad \text{and} \quad \text{sopp}(\rho_\mu) \subset \left] -\frac{1}{\mu}, \frac{1}{\mu} \right[. \quad (2.34)$$

$$\int_{\mathbb{R}} \rho_\mu(s) ds = 1 \quad \forall \mu \in \mathbb{N} \quad (2.35)$$

Let us denote by  $(f_{\nu\mu})_{\mu \in \mathbb{N}}$  the sequence of bounded function defined by:

$$f_{\nu\mu} = f_\nu * \rho_\mu \text{ for a fixed } \nu.$$

Then we have  $f_{\nu\mu}$  is a  $C^\infty$  bounded function we now consider the following approximated problem.

$$\begin{cases} u''_{\nu\mu} - \Delta u_{\nu\mu} + f(u_{\nu\mu}) = V & \text{in } Q; \\ u_{\nu\mu}(0) = u_0, u'_{\nu\mu} = u_1 & \text{in } \Omega; \\ u_{\nu\mu}(x, t) = 0 & \text{on } \Sigma. \end{cases} \quad (2.36)$$

As well known that for every  $(u_0, u_1, V) \in H_0^1(\Omega) \times L^2(\Omega) \times L^2(Q)$  there exist an unique solution for system (2.33). In order to obtain the existence of a solution for the system (1.3), let us suppose that,  $V, u_0, u_1$  be test function, then we have that:

$$u_{\nu\mu} \in L^\infty(0, T; H_0^1(\Omega) \cap H^2(\Omega)). \quad (2.37)$$

$$u_{\nu\mu} \in L^\infty(0, T; H_0^1(\Omega)) \quad (2.38)$$



$$u_{\nu\mu} \in L^\infty(0, T; H^2(\Omega)) \quad (2.39)$$

From Remark (2.4) we have that the normal derivate of  $u_{\nu\mu}$ , satisfies the following inequality:

$$\begin{aligned} \frac{1}{2} \int_{\Sigma} \left| \frac{\partial u_{\nu\mu}}{\partial \eta} \right|^2 d\Sigma &\leq C(n+1)(2+T) \sup J_{\nu\mu}(t) + \\ &+ cn \left\{ \frac{1}{2} |V|_Q^2 + \int_Q |G_{\nu\mu}(u_{\nu\mu})| \right\} dx dt. \end{aligned} \quad (2.40)$$

Where by  $J_{\nu\mu}(t)$  we are denoting the quadratic term associated to system (2.36), that is:

$$J_{\nu\mu}(t) = \frac{1}{2} |u_{\nu\mu}(t)|_\Omega^2 + \frac{1}{2} \|u_{\nu\mu}(t)\|^2$$

By Remarks 2.5, we have that,

$$\begin{aligned} J_{\nu\mu}(t) + \int_\Omega G_{\nu\mu}(u_{\nu\mu}) + b |u_{\nu\mu}|^2 dx &\leq \frac{1}{2} |V|_Q^2 + \\ &+ 2\mathcal{C}_Q E_{\nu\mu}(0) + \mathcal{C}_0 \int_0^T J_{\nu\mu}(s) ds. \end{aligned} \quad (2.41)$$

and since  $b$  is a positive number, and  $G_{\nu\mu}$  is uniformly bounded for all  $\mu \in \mathbb{N}$ , and

a fixed  $\nu$ , we have by Gronwall inequality that there exist a constant  $C_\nu$  such us:

$$J_{\nu\mu}(t) = \frac{1}{2} |u'_{\nu\mu}(t)|_\Omega^2 + \frac{1}{2} \|u_{\nu\mu}(t)\|^2 \leq C_\nu. \quad (2.42)$$

where  $E_{\nu\mu}(t)$  is the energy associated to system (2.36) with non quadratic term  $G_{\nu\mu}(u_{\nu\mu})$ .

Finally from Remark 2.6 we obtain that:

$$\begin{aligned} \int_Q u_{\nu\mu} f_{V\mu}(u_{\nu\mu}) dx \leq + \\ + \frac{1}{2} |V|_Q^2 + (3T + 2C) \sup_{[0,T]} J_{V\mu}(t). \end{aligned} \quad (2.43)$$

Relation (2.41), (2.42), and (2.43) are valid when  $V$ ,  $u_0$  and  $u_1$  are test function.

If we take a sequence of test function  $(V_m, u_{0m}, u_{1m})$  satisfying,

$$(V_m, u_{0m}, u_{1m}) \rightarrow (V, u_0, u_1) \text{ strongly in } L^2(Q) \times H_0^1(\Omega) \times L^2(\Omega)$$

certainly we have that the corresponding solutions  $u_{\nu u_m}$  converge to  $u_{\nu u}$  solution of system (2.36), when the datas  $V$ ,  $u_0$  and  $u_1$  are  $L^2(Q)$ ,  $H_0^1(\Omega)$  and  $L^2(\Omega)$

respectively. Moreover we have:

$$u_{\nu u_m} \rightarrow u_{\nu u} \text{ strongly in } L^\infty(O, T; H_0^1(\Omega)). \quad (2.44)$$

$$u_{\nu u_m} \rightarrow u_{\nu u} \text{ strongly in } L^\infty(O, T; L^2(\Omega)) \quad (2.45)$$

From (2.44) and (2.45) we conclude that relations (2.40), (2.41), (2.42) and (2.43) are valid when  $(V, u_0, u_1)$  belongs to  $L^2(Q) \times H_0^1(\Omega) \times L^2(\Omega)$  and  $u_{\nu u}$  is solution of (3.36).

On the other hand, by (2.42) we obtain that there exists a subsequence of  $(u_{\nu\mu})_{\mu \in \mathbb{N}}$  satisfying:

$$u'_{\nu\mu} \rightarrow u_\nu \text{ weak star in } L^\infty(O, T; H_0^1(\Omega)) \quad (2.46)$$

$$u'_{\nu\mu} \rightarrow u'_\nu \text{ weak star in } L^\infty(O, T; L^2(\Omega)) \quad (2.47)$$

from (2.46) and (2.47) we have that there exist another subsequence (that we still denoting in the same way) such that,

$$u_{\nu\mu} \rightarrow u_\nu \text{ strongly in } L^2(Q). \quad (2.48)$$

$$u_{\nu\mu} \rightarrow u_\nu \text{ a. e. in } Q. \quad (2.49)$$

$$f_{\nu\mu}(u_{\nu\mu}) \rightarrow f_\nu(u_\nu) \text{ a. e. in } Q. \quad (2.50)$$

$$G_{\nu\mu}(u_{\nu\mu}) \rightarrow G_\nu(u_\nu) \text{ a. e. in } Q. \quad (2.51)$$

Since  $f_{\nu\mu}$  is bounded for all  $\mu \in \mathbb{N}$  ( $\nu$  fixed), then  $G_{\nu\mu}$  is a Lipschitz's in  $\mathbb{R}$ , then by Lebesgue dominated convergence theorem we conclude that:

$$f_{\nu\mu}(u_{\nu\mu}) \rightarrow f_\nu(u_\nu) \text{ strongly in } L^2(Q). \quad (2.52)$$

$$G_{\nu\mu}(u_{\nu\mu}) \rightarrow G_\nu(u_\nu) \text{ strongly in } L^2(Q). \quad (2.53)$$

Now, from (2.48) and (2.52) we obtain:

$$u_{\nu\mu} \rightarrow u_\nu \text{ strongly in } L^\infty(O, T; H_0^1(\Omega)). \quad (2.54)$$

$$u'_{\nu\mu} \rightarrow u'_\nu \text{ strongly in } L^\infty(O, T; L^2(\Omega)). \quad (2.55)$$

Then by (2.40), (2.42) and (2.53) we obtain that there exist a subsequence of  $\frac{\partial u}{\partial \eta} \nu_\mu$ , which we still denote in the same way and a element  $\mathcal{X}$  in  $L^2(\Sigma)$  such

that:

$$\frac{\partial u_{\nu\mu}}{\partial \eta} \rightarrow \mathcal{X}_\nu \text{ weak in } L^2(\Sigma) \quad (2.56)$$

But since:

$$\frac{\partial u_{\nu\mu}}{\partial \eta} \rightarrow \frac{\partial u_\nu}{\partial \eta} \text{ weak in } H^{-1}(O, T; H^{1/2}(\Gamma)).$$

we conclude that  $\mathcal{X}_\nu = \frac{\partial u_\nu}{\partial \eta}$ . From (2.54) and (2.55) we have in particular that:

$$J_{\nu\mu}(t) \rightarrow J_\nu(t) \text{ uniformly on } [O, T]. \quad (2.57)$$

$$E_{\nu\mu}(t) \rightarrow E_\nu(t) \text{ uniformly on } [O, T]. \quad (2.58)$$

Then from (2.40), (2.53) and (2.57) we obtain that:

$$\begin{aligned} \frac{1}{2} \int_{\Sigma} \left| \frac{\partial u_\nu}{\partial \eta} \right|^2 d\Sigma &\leq C(n+1)(2+T) \sup_{[O, T]} J_\nu(t) + \\ &+ Cn \left\{ \frac{1}{2} |V|_Q^2 + \int_Q |G_\nu(u_\nu)| \right\} dx dt. \end{aligned} \quad (2.59)$$

Now by (2.41), (2.43), (2.48), (2.53), (2.57) and (2.58) we obtain:

$$\begin{aligned} J_\nu(t) + \int_{\Omega} \{G_\nu(u_\nu) + b|u_\nu|^2\} &\leq \\ \leq \frac{1}{2} |V|_Q^2 + 2C_0 E_\nu(0) + C_0 \int_0^t J_\nu(s) ds. \end{aligned} \quad (2.60)$$

$$\int u_\nu f_\nu(u_\nu) dx dt \leq \frac{1}{2} |V|_Q^2 + (3T + 2C) \sup_{[0, T]} J_\nu(t). \quad (2.61)$$

From (2.10) and (2.59) taking  $b$  such that  $G_\nu(u_\nu) + b|u_\nu|^2$  be positive, we obtain:

$$J_\nu(t) \leq \frac{1}{2} |V|_Q^2 + 2C_0 E_\nu(0) + C_0 \int_0^t J_\nu(s) ds. \quad (2.62)$$

Now by Gronwal's inequality we obtain that:

$$J_\nu(t) \leq \left\{ \frac{1}{2} |V|_Q^2 + 2C_0 E_\nu(0) \right\} e^{C_0 T}, \forall t \in [0, T]. \quad (2.63)$$

Since,

$$E_\nu(0) = \frac{1}{2} \left\{ \|u_0\|^2 + |u_1|_\Omega^2 \right\} + \int_\Sigma G_\nu(u_0) dx,$$

we conclude from (2.12) and the hypothesis of Theorem 2.6, that the second member of (2.63) is bounded by a constand  $C_1 > 0$ , independing of  $\nu$ , then from (2.63) we have:

$$\sup_{[0, T]} J_\nu(t) \leq \left\{ \frac{1}{2} |V|_Q^2 + 2C_0 E_\nu(0) \right\} e^{C_0 T} \leq C_1. \quad (2.64)$$

Then we have that there exist a subsequence of  $(u_\eta)_{\eta \in \mathbb{N}}$ , that we still

denote on the same way, and a element  $u \in L^\infty(O, T; H_0^1(\Omega))$  such that  $u' \in L^\infty(O, T; L^2(\Omega))$ , satisfying:

$$u \in L^\infty(O, T; H_0^1(\Omega)) \text{ such that } u' \in L^\infty(O, T; L^2(\Omega)), \quad (2.65)$$

$$u'_\nu \rightarrow u' \text{ weak star in } L^\infty(O, T; L^2(\Omega)). \quad (2.66)$$

By (2.61) and (2.63) we obtain that:

$$\int_Q u_\nu f_\nu(u_\nu) dx \leq 3C_2 e^{C_0 T} \{ |V|_Q^2 + E_\nu(0) \} \leq C_3. \quad (2.67)$$

where  $C_2 = \max\{n, T, C, C_0\}$ . But from (2.9) we obtain that  $|u_\nu f_\nu(u_\nu)| \leq u_\nu f_\nu(u_\nu) + 2C_0(u_\nu^2 + 1)$ , from where we have:

$$\int_Q |u_\nu f_\nu(u_\nu)| dx \leq C_3 + 4C_0 x C_1 \text{ med } (Q) = C_4. \quad (2.68)$$

Then by (2.68) and from Theorem 1.1 of W.A. Strauss [4] we have that:

$$f_\nu(u_\nu) \rightarrow f_{(u)} \text{ strongly in } L^1(Q). \quad (2.69)$$

Then we conclude that  $u$  is a solution of problem (1.1). Finally from (2.59)

and (2.63) we have that there exist a constant  $\mathcal{C}_5$  (independent of  $\nu$ ) such that:

$$\int_{\Sigma} \left| \frac{\partial u_\nu}{\partial \eta} \right| d\Sigma \leq \mathcal{C}_5 \left\{ |V|_Q^2 + E_\nu(0) + \int_Q |G_\nu(u_\nu)| dx dt \right\}. \quad (2.70)$$

But from (2.60) and (2.64) we have that exist a constant  $\mathcal{C}_6$  such that,

$$\int_Q G_\nu(u_\nu) dx \leq \mathcal{C}_6 \left\{ |V|_Q^2 + E_\nu(0) \right\}. \quad (2.71)$$

Now from (2.10) we have that  $G_\nu(t) \leq -\mathcal{C}_0(t^2 + 2|t|)$ , from where we conclude that  $|G_\nu(t) \leq G_\nu(t)| + 2\mathcal{C}_0(t^2 + 2|t|)$  taking  $t = u_n$  we have after to integrate in  $Q$ , that:

$$\int_Q G_\nu(u_\nu) dx dt \leq \int_Q G_\nu(u_\nu) dx dt + 2\mathcal{C}_0 \int_Q u_\nu^2 + 2|u_\nu| dx dt. \quad (2.72)$$

On the other hand, there exist a constant  $\mathcal{C}_7$  such that:

$$\int_Q |u_\nu| + 2|u_\nu| dx dt \leq \mathcal{C}_7 \sup_{[0,T]} J_\nu(t). \quad (2.73)$$

From (2.70), (2.71), (2.72) and (2.73) we obtain another constant, say  $\mathcal{C}_8$  such



that:

$$\int_{\Sigma} \left| \frac{\partial u_{\nu}}{\partial \eta} \right| d\Sigma \leq C_8 \{ |V|_Q^2 + E_{\nu}(0) \} \quad (2.74)$$

Since the second member of (2.74) is bounded we obtain a subsequence of,

$$\left( \frac{\partial u_{\nu}}{\partial \eta} \right)_{\nu \in \mathbb{N}}$$

and a element  $\mathcal{X}$  in  $L^2(\Sigma)$  such that:

$$\frac{\partial u_{\nu}}{\partial \eta} \rightarrow \mathcal{X} \text{ weak in } L^2(\Sigma). \quad (2.75)$$

But,

$$\frac{\partial u_{\nu}}{\partial \eta} \rightarrow \frac{\partial u}{\partial \eta} \text{ in } H^{-1}(0, T; W^{-1/p-2, p'}(\Gamma)),$$

Where  $p > \frac{n}{2}$ , then we have that  $\mathcal{X} = \frac{\partial u}{\partial \eta}$ , and letting  $\nu \rightarrow \infty$  in (2.74) we have,

$$\int_{\Sigma} \left| \frac{\partial u}{\partial \eta} \right|^2 d\Sigma \leq C_8 \{ |V|_Q^2 + E(0) \} \quad \text{Q.E.D.}$$

## REFERENCES

- 1.- J. L. Lions - Hidden regularity in some nonlinear hyperbolic equations.

*Matemática Aplicada e Computacional*, v.6, n.1, pp. 7-15, (1987).

- 2.- J. L. Lions - Quelques méthodes de résolution des problèmes aux limites nonlineares, Paris. Dunod, (1969).
- 3.- J. L. Lions - Controlabilité exacte de systèmes distribués, vol. 1, Collège de France, Département de Mathématiques, n. 3., Rue d'Ulm 75005, Paris, France.
- 4.- W. A. Strauss - On weak solution of semilinear hyperbolic equations, *An Acad. Brasil. Ciências*, vol. 42, n,4, pp. 645 - 651, (1970).