

N-DIMENSIONS LINEAR THERMOELASTIC SYSTEM IN DOMAINS WITH MOVING BOUNDARY

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ABSTRACT

In this article the authors investigated the existence of solutions for the linear thermoelastic system in a noncylindrical domain \hat{Q} of \mathfrak{R}^{n+1} .

1. Introduction

Let $K: [0, T] \rightarrow \mathfrak{R}$ be a C^2 function and let Ω be an open bounded set in \mathfrak{R}^n , with regular boundary. Let us consider the family of open sets in \mathfrak{R}^{n+1} :

$$\Omega_t = \{K(t)y; y \in \Omega\}, \quad 0 \leq t \leq T,$$

with lateral boundary Γ_t . We define the noncylindrical domain \hat{Q} in \mathfrak{R}^{n+1} by:

$$Q = \bigcup_{0 < t < T} \Omega_t \times \{t\},$$

with lateral boundary

$$\Sigma = \bigcup_{0 < t < T} \Gamma_t \times \{t\}.$$

In this work we investigate the existence of solutions for the following thermoelastic system in \hat{Q} :

$$\left\{ \begin{array}{l} u'' - \mu \Delta u - (\lambda + \mu) \nabla \operatorname{div} u + \alpha \nabla \theta = 0 \quad \text{in } \hat{Q} \\ \theta' - \Delta \theta + \beta \operatorname{div} u_t = 0 \quad \text{in } \hat{Q} \\ u = 0, \theta = 0 \quad \text{in } \hat{\Sigma} \\ u(x, 0) = u_0(x), u'(x, 0) = u_1(x), \theta(x, 0) = \theta_0(x) \quad \text{in } \Omega_0. \end{array} \right. \quad (*)$$

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We denote by $x = (x_1, \dots, x_n)$ a point in \mathfrak{R}^n while t stands for the time variable. The displacement-vector is denoted by $u(x, t) = (u_1(x, t), \dots, u_n(x, t))$ and the thermal difference by $\theta = \theta(x, t)$; $\lambda, \mu > 0$ are Lemé's constants and $\alpha, \beta > 0$ the coupling parameters.

There is an extensive literature for the study of (*) in cylindrical domains. It can be seen in [1], [3], [4], [5] and [6]. For the one-dimensional noncylindrical case, in [2], it is studied the existence, uniqueness and energy decay of weak solutions. We don't know results for (*) in n-dimensional case in noncylindrical domains.

2. Notations and main Results

Let us consider the real function $K(t)$ satisfying to the following conditions:

H1) $K \in C^3([0, T]; \mathfrak{R})$ with $K_0 = \min_{0 \leq t \leq T} K(t) > 0$.

H2) There exists a positive constant K_1 , such that

$$\left[\mu \delta_i^j - (K')^2 y_i y_j \right] \epsilon_i \epsilon_j \geq K_1 |\epsilon|^2_{\mathfrak{R}^n} \quad \forall \epsilon \in \mathfrak{R}^n, \forall y \in \Omega, \quad 0 \leq t \leq T.$$

Remark. It is noted that when (x, t) varies in \hat{Q} the point (y, t) with $y = K^{-1}(t)x$ varies in the cylinder $Q = \Omega \times]0, T[$. Thus, the application

$$\tau: \hat{Q} \rightarrow Q$$

given by $\tau(x, t) = (y, t)$, is a diffeomorfism. Let us consider the change of variable in (*), given by

$$\begin{aligned} v(y, t) &= u(K(t)y, t) \\ \phi(y, t) &= \theta(K(t)y, t), \end{aligned}$$

We obtain the following mixed problem:

$$\left\{ \begin{aligned} &v'' - \frac{\partial}{\partial y_j} \left[a_{ij}(y, t) \frac{\partial v}{\partial y_i} \right] - \frac{(\lambda + \mu)}{K^2} \nabla \operatorname{div} v + \frac{\alpha}{K} \nabla \phi + b_i(y, t) \frac{\partial v}{\partial y_i} \\ &+ c_i(y, t) \frac{\partial v}{\partial y_i} = 0 \quad \text{in } Q \\ &\phi' - \frac{1}{K^2} \Delta \phi + \frac{\beta}{K} \operatorname{div} v' + d_i(y, t) \frac{\partial \phi}{\partial y_i} + \frac{\beta K'}{K^2} \operatorname{div} v \\ &+ e_i(y, t) \frac{\partial^2 v_j}{\partial y_j \partial y_i} = 0 \quad \text{in } Q \\ &v = 0, \phi = 0 \quad \text{in } \Sigma \\ &v(y, 0) = v_0(y), v'(y, 0) = v_1(y), \phi(y, 0) = \phi_0(y), \quad y \in \Omega \end{aligned} \right. \quad (**)$$

where

$$a_{ij}(y, t) = [\mu \delta_i^j - (K')^2 y_i y_j] K^{-2}, \quad b_i(y, t) = \frac{-2K'}{K} y_i$$

$$c_i(y, t) = [(1-n)(K')^2 - K'' K] K^{-2} y_i, \quad d_i(y, t) = \frac{-y_i K'}{K}, \quad e_i(y, t) = \frac{-\beta y_i K'}{K^2}.$$

Let us represent by $((\cdot), \|\cdot\|)$ and $(\cdot, |\cdot|)$, the scalar product and norm in $[H_0^1(\Omega)]^n$ and $[L^2(\Omega)]^n$ respectively. We denote by $a(t, v, w)$ the bilinear form defined in $[H_0^1(\Omega)]^n$ by:

$$a(t, v, w) = \int_{\Omega} a_{ij}(y, t) \frac{\partial v}{\partial y_i} \cdot \frac{\partial w}{\partial y_j} dy.$$

Note that $a(t, v, w)$ is continuous, symmetric and coercive.

The main results in our work are:

Teorema 1. Given $u_0 \in [H_0^1(\Omega_0) \cap H^2(\Omega_0)]^n$, $\theta_0 \in H_0^1(\Omega_0) \cap H^2(\Omega_0)$, and $u_1 \in [H_0^1(\Omega_0)]^n$, then there exist functions

$$u : \hat{Q} \rightarrow \mathfrak{R}^n \text{ and } \theta : \hat{Q} \rightarrow \mathfrak{R}$$

such that:

$$u \in L^\infty(0, T; [H_0^1(\Omega_t) \cap H^2(\Omega_t)]^n), u' \in L^\infty(0, T; [H_0^1(\Omega_t)]^n), u'' \in L^\infty(0, T; [L^2(\Omega_t)]^n)$$

$$\theta \in L^2(0, T; H_0^1(\Omega_t) \cap H^2(\Omega_t)), \theta' \in L^2(0, T; H_0^1(\Omega_t)),$$

and they are solutions of the problem (*) in \hat{Q} .

Teorema 2. Given $v_0 \in [H_0^1(\Omega) \cap H^2(\Omega)]^n$, $\phi_0 \in H_0^1(\Omega) \cap H^2(\Omega)$, and $v_1 \in [H_0^1(\Omega)]^n$, then there exist functions

$$v : Q \rightarrow \mathfrak{R}^n \text{ and } \phi : Q \rightarrow \mathfrak{R}$$

such that:

$$v \in L^\infty(0, T; [H_0^1(\Omega) \cap H^2(\Omega)]^n), v' \in L^\infty(0, T; [H_0^1(\Omega)]^n), v'' \in L^\infty(0, T; [L^2(\Omega)]^n)$$

$$\phi \in L^2(0, T; H_0^1(\Omega) \cap H^2(\Omega)), \phi' \in L^2(0, T; H_0^1(\Omega)),$$

and they are solutions of the problem (**) in Q .

Proof of theorem 2. We consider $\{w_j\}$, $j=1,2,\dots$, a Hilbertien basis in $[H_0^1(\Omega) \cap H^2(\Omega)]^n$.

We represent by $V_m = [w_1, \dots, w_m]$ the subspace generate by the vectors w_1, \dots, w_m . We

find $v_m(t)$ and $\phi_m(t)$ in V_m , solutions of the following system of ordinary differential equations:

$$\left(v_m'', w \right) + a(t, v_m, w) - \left(\frac{(\lambda + \mu)}{K^2} \nabla \operatorname{div} v_m, w \right) + \frac{\alpha}{K} (\nabla \phi_m, w) + \left(b_i \frac{\partial v_m'}{\partial y_i}, w \right) \quad (1)$$

$$+ \left(c_i \frac{\partial v_m}{\partial y_i}, w \right) = 0 \quad \forall w \in V_m.$$

$$\left(\phi_m', w \right) + \frac{1}{K^2} (\nabla \phi_m, \nabla w) + \left(\frac{\beta}{K} \operatorname{div} v_m', w \right) + \left(d_i \frac{\partial \phi_m}{\partial y_i}, w \right) \quad (2)$$

$$+ \left(\frac{\beta K'}{K^2} \operatorname{div} v_m, w \right) + \left(e_i \frac{\partial^2 v_{jm}}{\partial y_j \partial y_i}, w \right) = 0, \quad \forall w \in V_m.$$

$$v_m(0) = v_{0m} \rightarrow v_0 \text{ in } [H_0^1(\Omega) \cap H^2(\Omega)]^n$$

$$v_m'(0) = v_{1m} \rightarrow v_1 \text{ in } [H_0^1(\Omega)]^n$$

$$\phi_m(0) = \phi_{0m} \rightarrow \phi_0 \text{ in } H_0^1(\Omega) \cap H^2(\Omega).$$

First Estimate

Taking $w = v_m'$ in (1) and $w = \phi_m$ in (2), we have:

$$\frac{1}{2} \frac{d}{dt} |v_m'|^2 + \frac{1}{2} \frac{d}{dt} a(t, v_m, v_m) - \frac{1}{2} a'(t, v_m, v_m) - \left(\frac{(\lambda + \mu)}{K^2} \nabla \operatorname{div} v_m, v_m' \right) \quad (3)$$

$$+ \frac{\alpha}{K} (\nabla \phi_m, v_m') + \left(b_i \frac{\partial v_m'}{\partial y_i}, v_m' \right) + \left(c_i \frac{\partial v_m}{\partial y_i}, v_m' \right) = 0.$$

$$\frac{1}{2} \frac{d}{dt} |\phi_m|^2 + \frac{1}{K^2} \|\phi_m\|^2 + \left(\frac{\beta}{K} \operatorname{div} v_m', \phi_m \right) + \left(d_i \frac{\partial \phi_m}{\partial y_i}, \phi_m \right) \quad (4)$$

$$+ \left(\frac{\beta K'}{K^2} \operatorname{div} v_m, \phi_m \right) + \left(e_i \frac{\partial^2 v_{jm}}{\partial y_j \partial y_i}, \phi_m \right) = 0.$$

Developing some terms of (3) and (4), we obtain:

$$-\left(\frac{(\lambda + \mu)}{K^2} \nabla \operatorname{div} v_m, v'_m\right) = \frac{(\lambda + \mu)}{K^2} (\operatorname{div} v_m, \operatorname{div} v'_m) \quad (5)$$

$$\begin{aligned} &= \frac{1}{2} \frac{d}{dt} \left(\frac{(\lambda + \mu)}{K^2} |\operatorname{div} v_m|^2 \right) + \frac{(\lambda + \mu)K'}{K^3} |\operatorname{div} v_m|^2 \\ &\leq \frac{1}{2} \frac{d}{dt} \left(\frac{(\lambda + \mu)}{K^2} |\operatorname{div} v_m|^2 \right) + c_1 \|v_m\|^2 \end{aligned}$$

$$\frac{\alpha}{K} (\nabla \phi_m, v'_m) = \frac{\alpha}{K} \int_{\Omega} \nabla \phi_m v'_m dy \quad (6)$$

$$-\frac{\alpha}{K} \int_{\Omega} \phi_m \operatorname{div} v'_m dy = -\left(\frac{\alpha}{K} \operatorname{div} v'_m, \phi_m\right)$$

$$|a'(t, v_m, v'_m)| \leq c_2 \|v_m\|^2 \quad (7)$$

$$\left| \left(b_i \frac{\partial v'_m}{\partial y_i}, v'_m \right) \right| = \left| \frac{1}{2} \int_{\Omega} b_i \frac{\partial |v'_m|^2}{\partial y_i} dy \right|$$

$$= \frac{1}{2} \left| - \int_{\Omega} \frac{\partial b_i}{\partial y_i} |v'_m|^2 dy + \int_{\Gamma} b_i |v'_m|^2 \eta_i d\Gamma \right| \quad (8)$$

$$\leq c_3 |v'_m|^2$$

$$\left| \left(c_i \frac{\partial v_m}{\partial y_i}, v'_m \right) \right| \leq c_4 (\|v_m\|^2 + |v'_m|^2) \quad (9)$$

$$\left| \left(d_i \frac{\partial \phi_m}{\partial y_i}, \phi_m \right) \right| = \frac{1}{2} \left| \int_{\Omega} d_i \frac{\partial |\phi_m|^2}{\partial y_i} dy \right|$$

$$= \left| - \frac{1}{2} \int_{\Omega} \frac{\partial d_i}{\partial y_i} |\phi_m|^2 dy + \frac{1}{2} \int_{\Gamma} |\phi_m|^2 d_i \eta_i d\Gamma \right| \leq c_5 |\phi_m|^2 \quad (10)$$

$$\left| \left(\frac{\beta K'}{K^2} \operatorname{div} v_m, \phi_m \right) \right| \leq c_6 \left(\|v_m\|^2 + |\phi_m|^2 \right) \quad (11)$$

$$\begin{aligned} \left(e_i \frac{\partial^2 v_{mj}}{\partial y_j \partial y_i}, \phi_m \right) &= \int_{\Omega} \phi_m \vec{e} \cdot \nabla (\operatorname{div} v_m) dy \\ &= - \int_{\Omega} \operatorname{div} \left(\phi_m \vec{e} \right) \cdot \operatorname{div} v_m dy + \int_{\Gamma} \phi_m \vec{e} \cdot \operatorname{div} v_m d\Gamma \end{aligned} \quad (12)$$

$$= - \int_{\Omega} \operatorname{div} \left(\phi_m \vec{e} \right) \cdot \operatorname{div} v_m dy \text{ where } \vec{e} = (e_1, \dots, e_n). \text{ Thus,}$$

$$\left| \left(e_i \frac{\partial^2 v_{mj}}{\partial y_j \partial y_i}, \phi_m \right) \right| \leq \frac{1}{2K^2} \|\phi_m\|^2.$$

From (3) – (12), we have:

$$\begin{aligned} &\frac{1}{2} \frac{d}{dt} \left\{ \frac{\beta}{\alpha} |v'_m|^2 + \frac{\beta}{\alpha} a(t, v_m, v_m) + \frac{\beta}{\alpha} \left(\frac{\lambda + \mu}{K^2} \right) \operatorname{div} |v_m|^2 + |\phi_m|^2 \right\} \\ &+ \frac{1}{2K^2} \|\phi_m\|^2 \leq c_8 \left(\|v_m\|^2 + |v'_m|^2 + |\phi_m|^2 \right) \end{aligned} \quad (13)$$

Integrating (13) on $[0, t[$ and applying Gronwall inequality, we conclude that there is a positive constant C independent of m and t , such that

$$\begin{aligned} &\frac{\beta}{\alpha} |v'_m|^2 + \frac{\beta}{\alpha} a(t, v_m, v_m) + \frac{\beta}{\alpha} \left(\frac{\lambda + \mu}{K^2} \right) \operatorname{div} |v_m|^2 \\ &+ |\phi_m|^2 + \frac{1}{2K^2} \int_0^t \|\phi_m\|^2 dt \leq C, \end{aligned} \quad (14)$$

Second Estimate

Differentiating (1) with respect to t and taking $w = v_m''$, we obtain:

$$\begin{aligned} & \left(v_m''', v_m'' \right) + a(t, v_m, v_m'') + a'(t, v_m, v_m'') - \left(\frac{\lambda + \mu}{K^2} \nabla \operatorname{div} v_m, v_m'' \right) \\ & + \frac{2(\lambda + \mu)K'}{K^3} \left(\nabla \operatorname{div} v_m, v_m'' \right) + \frac{\alpha}{K} \left(\nabla \phi_m, v_m'' \right) \\ & + \left(b_i \frac{\partial v_m''}{\partial y_i}, v_m'' \right) + \left(b_i' \frac{\partial v_m''}{\partial y_i}, v_m'' \right) + \left(c_i \frac{\partial v_m'}{\partial y_i}, v_m'' \right) + \left(c_i' \frac{\partial v_m''}{\partial y_i}, v_m'' \right) = 0 \end{aligned} \quad (15)$$

Differentiating (2) with respect to t and taking $w = \phi_m'$, we have:

$$\begin{aligned} & \left(\phi_m'', \phi_m' \right) + \frac{1}{K^2} \left| \nabla \phi_m' \right|^2 - \frac{2K'}{K^3} \left(\nabla \phi_m, \nabla \phi_m' \right) + \left(\frac{\beta}{K} \operatorname{div} v_m'', \phi_m' \right) \\ & + \left(d_i \frac{\partial \phi_m'}{\partial y_i}, \phi_m' \right) + \left(d_i' \frac{\partial \phi_m'}{\partial y_i}, \phi_m' \right) + \left(\left(\frac{\beta K'}{K^2} \right)' \operatorname{div} v_m, \phi_m' \right) \\ & + \left(e_i \frac{\partial^2 v_{jm}}{\partial y_i \partial y_i}, \phi_m' \right) + \left(e_i' \frac{\partial^2 v_{jm}}{\partial y_i \partial y_i}, \phi_m' \right) = 0. \end{aligned} \quad (16)$$

Developing term terms of (15) and (16), we get:

$$a'(t, v_m, v_m'') = \int_{\Omega} a_{ij}' \frac{\partial v_m}{\partial y_i} \frac{\partial v_m''}{\partial y_j} dy \quad (17)$$

$$= \frac{d}{dt} \int_{\Omega} a_{ij}' \frac{\partial v_m}{\partial y_i} \frac{\partial v_m'}{\partial y_j} dy - \int_{\Omega} a_{ij}'' \frac{\partial v_m}{\partial y_i} \frac{\partial v_m'}{\partial y_j} dy - \int_{\Omega} a_{ij}' \frac{\partial v_m'}{\partial y_i} \frac{\partial v_m'}{\partial y_j} dy \quad (18)$$

$$\begin{aligned} \frac{2(\lambda + \mu)K'}{K^3} \left(\nabla \operatorname{div} v_m, v_m'' \right) &= \frac{2(\lambda + \mu)K'}{K^3} \int_{\Omega} \operatorname{div} v_m \operatorname{div} v_m'' dy \\ &- \int_{\Omega} \frac{2(\lambda + \mu)K'}{K^3} \left| \operatorname{div} v_m' \right|^2 dy. \end{aligned}$$

$$\frac{-2K'}{K^3} (\nabla \phi_m, \nabla \phi'_m) = \frac{d}{dt} \left(\frac{-K'}{K^3} \|\phi_m\|^2 \right) + \left(\frac{K'}{K^3} \right) \|\phi_m\|^2 \quad (19)$$

$$\frac{-\alpha K'}{K^2} \int_{\Omega} \nabla \phi_m \nu_m'' dy \leq C_9 \left| \nu_m'' \right|^2 + \frac{1}{4K^2} \|\phi_m\|^2 \quad (20)$$

$$\begin{aligned} \left(e'_i \frac{\partial^2 \nu_{jm}}{\partial y_j \partial y_i}, \phi'_m \right) &= \int_{\Omega} e'_i \frac{\partial}{\partial y_j} \left(\frac{\partial \nu_{jm}}{\partial y_i} \right) \phi'_m dy = \int_{\Omega} e'_i \operatorname{div} \frac{\partial \nu_m}{\partial y_i} \phi'_m dy \\ &= - \int_{\Omega} \frac{\partial}{\partial y_i} (e'_i \phi'_m) \operatorname{div} \nu_m dy + \int_{\Gamma} \operatorname{div} \nu_m e'_i \phi'_m \eta_i d\Gamma \\ &= - \int_{\Omega} \frac{\partial e'_i}{\partial y_i} \phi'_m \operatorname{div} \nu_m dy - \int_{\Omega} e'_i \frac{\partial \phi'_m}{\partial y_i} \operatorname{div} \nu_m dy \\ &\leq c_{10} \left(\left| \phi'_m \right|^2 + \|\nu_m\|^2 \right) + \frac{1}{4K^2} \left\| \phi'_m \right\|^2. \end{aligned} \quad (21)$$

From (15)-(21), we obtain:

$$\begin{aligned} &\frac{1}{2} \frac{d}{dt} \left\{ \frac{\beta}{\alpha} \left| \nu_m'' \right|^2 + \frac{\beta}{\alpha} a(t, \nu'_m, \nu'_m) + \frac{\beta (\lambda + \mu)}{\alpha K^2} \left| \operatorname{div} \nu'_m \right|^2 + \left| \phi'_m \right|^2 \right. \\ &+ \left. \frac{\beta}{\alpha} \int_{\Omega} a'_{ij} \frac{\partial \nu_m}{\partial y_i} \frac{\partial \nu_m}{\partial y_j} dy + \frac{\beta}{\alpha} \int_{\Omega} \frac{2(\lambda + \mu)K'}{K^3} \operatorname{div} \nu_m \operatorname{div} \nu'_m dy - \frac{K'}{K^3} \|\phi_m\|^2 \right\} \\ &+ \frac{1}{4K^2} \left\| \phi'_m \right\|^2 \leq c \left[\left\| \nu'_m \right\|^2 + \left| \nu_m'' \right|^2 + \left| \phi'_m \right|^2 + \|\nu_m\|^2 + \|\phi_m\|^2 \right] \end{aligned} \quad (22)$$

Integrating (22) on $[0, t[$, and applying Gronwall inequality, we deduce:

$$\begin{aligned} &\frac{\beta}{\alpha} \left| \nu_m'' \right|^2 + \frac{\beta}{2\alpha} \rho_0 \left\| \nu'_m \right\|^2 + \frac{\beta (\lambda + \mu)}{\alpha K^2} \left| \operatorname{div} \nu'_m \right|^2 + \left| \phi'_m \right|^2 \\ &+ \int_0^t \frac{1}{8K^2} \left\| \phi'_m \right\|^2 dt \leq C \end{aligned} \quad (23)$$

From (14) and (23), it follows that:

$$\begin{cases}
 v_m'' \rightarrow v'' \text{ weak star in } L^\infty(0, T; [L^2(\Omega)]^n) \\
 v_m \rightarrow v \text{ weak star in } L^\infty(0, T; [H_0^1(\Omega)]^n) \\
 \frac{1}{K} \phi_m \rightarrow \frac{1}{K} \phi \text{ weak star in } L^\infty(0, T; H_0^1(\Omega)) \\
 \phi_m' \rightarrow \phi' \text{ weak star in } L^2(0, T; H_0^1(\Omega))
 \end{cases} \quad (24)$$

From (24) taking limits in (1) and (2), we have the proof of the theorem 2.

Proof of theorem 1.

We set

$$\begin{aligned}
 \theta_0(y) &= u_0(K(0)y), \quad \phi_0(y) = \theta_0(K(0)y), \\
 u_1(y) &= K'(0) \nabla u_0(K(0)y) \cdot y + u_1(K(0)y)
 \end{aligned} \quad (25)$$

that satisfy the hypothesis of theorem 2.

Therefore, there exist θ , ϕ solutions of the problem (**). Using θ and ϕ we construct u and v given by:

The regularity of u and v can be obtained from θ , ϕ and the hypothesis H1). We also verify the following identities:

$$\alpha \nabla \theta = \frac{\alpha}{K} \nabla \phi \quad (27)$$

$$(\lambda + \mu) \nabla \operatorname{div} u = \frac{\lambda + \mu}{K^2} \nabla \operatorname{div} v \quad (28)$$

$$\mu \Delta u = \frac{\mu}{K^2} \Delta v \quad (29)$$

$$u'' = v'' - \frac{\partial}{\partial y_j} \left(a_{ij} \frac{\partial v}{\partial y_i} \right) + \frac{\mu}{K^2} \Delta v + b_i \frac{\partial v}{\partial y_i} + c_i \frac{\partial v}{\partial y_i} \quad (30)$$

$$\Delta \theta = \frac{1}{K^2} \Delta \phi \quad (31)$$

$$\beta \operatorname{div} u' = \frac{\beta}{K} \operatorname{div} v' + \frac{\beta K'}{K^2} \operatorname{div} v + e_i \frac{\partial^2 v_j}{\partial y_j \partial y_i} \quad (32)$$

$$\theta' = \phi' + d_i \frac{\partial \phi}{\partial y_i} \quad (33)$$

from (27)-(33), using the fact that (v, ϕ) satisfies (**), we deduce that (u, θ) satisfies (*), and hence the theorem 1 is completely proved.

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